

Lecture Notes 15

Riemannian Geodesics

Here we show that every Riemannian manifold admits a unique connection, called the Riemannian or Levi-Civita connection, which satisfies two properties: symmetry, and compatibility with the metric, as we describe below. This result is known as the fundamental theorem of Riemannian geometry. Further we will show that the geodesics which arise from a Riemannian connection are locally minimize distance.

0.1 The bracket

For any pair of vector fields $X, Y \in \mathcal{X}(M)$ we may define a new vector field $[X, Y] \in \mathcal{X}(M)$ as follows. First recall that T_pM is isomorphic to D_pM the space of derivations of the germ of functions of M at p . Thus we may define $[X, Y]$ by describing how it acts on functions at each point:

$$[X, Y]_p f := X_p(Yf) - Y_p(Xf).$$

One may check that this does indeed define a derivation, i.e., $[X, Y]_p(\lambda f + g) = \lambda[X, Y]_p f + [X, Y]_p g$, and $[X, Y]_p(fg) = ([X, Y]_p f)g(p) + f(p)([X, Y]_p g)$. Further note that if $e_i(p) = e_i$ denotes the standard basis vector field of \mathbf{R}^n then $[e_i, e_j] = 0$ (since partial derivatives commute). On the other hand it is not difficult to construct examples of vector fields whose bracket does not vanish:

Example 0.1. Let X, Y be vector fields on \mathbf{R}^2 given by $X(x, y) = (1, 0)$ and $Y(x, y) = (0, x)$. Then

$$[X, Y]f = X\left(x \frac{\partial f}{\partial y}\right) - Y\left(\frac{\partial f}{\partial x}\right) = \frac{\partial f}{\partial y} + x \frac{\partial^2 f}{\partial x \partial y} - x \frac{\partial^2 f}{\partial y \partial x} = \frac{\partial f}{\partial y}$$

Lemma 0.2. Let $f: M \rightarrow N$ be a diffeomorphism, and $X, Y \in \mathcal{X}(M)$. Then

$$df([X, Y]) = [dfX, dfY].$$

Proof. Recall that for any vectorfield Z on M and function g on N , we have

$$((dfZ)g)(f(p)) = (dfZ)_{f(p)}g = (df_p Z)g = Z_p(g \circ f).$$

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Thus if we let $\bar{Z} := dfZ$, and $\bar{p} := f(p)$, then

$$\left((\bar{Z}g) \circ f \right) (p) = (\bar{Z}g)(\bar{p}) = \bar{Z}_{\bar{p}}g = Z_p(g \circ f) = \left(Z(g \circ f) \right) (p).$$

Using the last set of identities, we may now compute

$$\begin{aligned} \left(\overline{[X, Y]g} \right) (\bar{p}) &= [X, Y]_p(g \circ f) \\ &= X_p(Y(g \circ f)) - Y_p(X(g \circ f)) \\ &= X_p\left((\bar{Y}g) \circ f \right) - Y_p\left((\bar{X}g) \circ f \right) \\ &= \bar{X}_{\bar{p}}(\bar{Y}g) - \bar{Y}_{\bar{p}}(\bar{X}g) \\ &= \left(\overline{[X, Y]g} \right) (\bar{p}). \end{aligned}$$

□

Corollary 0.3. *Let (U, ϕ) be a local chart of M and $E_i(p) := d\phi_{\phi(p)}^{-1}(e_i)$ be the associated coordinate vector fields on U . Then $[E_i, E_j] = 0$.* □

Exercise 0.4. Show that the bracket satisfies the following properties:

$$[X, Y] = -[Y, X] \quad \text{and} \quad [X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0$$

0.2 Riemannian Connections

Recall that the standard connection in \mathbf{R}^n is defined as

$$\nabla_X Y := (X(Y^1), \dots, X(Y^n)).$$

Furthermore, recall that in \mathbf{R}^n , for any function $f: \mathbf{R}^n \rightarrow \mathbf{R}$ and vector field X we have

$$Xf = \langle X, \text{grad } f \rangle.$$

Using these identities, we compute that

$$\begin{aligned} Z\langle X, Y \rangle &= \sum \langle Z, \text{grad}(X^i Y^i) \rangle \\ &= \sum \langle Z, \text{grad}(X^i) Y^i + X^i \text{grad}(Y^i) \rangle \\ &= \sum \langle Z, \text{grad } X^i \rangle Y^i + \sum \langle Z, \text{grad } Y^i \rangle X^i \\ &= \langle \nabla_Z X, Y \rangle + \langle X, \nabla_Z Y \rangle. \end{aligned}$$

Motivated by this observation we say that a connection on a Riemannian manifold (M, g) is compatible with the metric provided that

$$Zg(X, Y) = g(\nabla_Z X, Y) + g(X, \nabla_Z Y).$$

Further note that

$$\langle Y, \nabla_X \text{grad } f \rangle - \langle X, \nabla_Y \text{grad } f \rangle = \sum_{i=1}^n Y^i X^i D_i f - \sum_{i=1}^n X^i Y^i D_i f = 0.$$

This property, together with the compatibility of ∇ with the innerproduct which we established above, may be used to compute that

$$\begin{aligned} (\nabla_X Y - \nabla_Y X)f &= \langle \nabla_X Y, \text{grad } f \rangle - \langle \nabla_Y X, \text{grad } f \rangle + \langle Y, \nabla_X \text{grad } f \rangle - \langle X, \nabla_Y \text{grad } f \rangle \\ &= X \langle Y, \text{grad } f \rangle - Y \langle X, \text{grad } f \rangle \\ &= X(Yf) - Y(Xf) \\ &= [X, Y](f). \end{aligned}$$

Thus we say that a connection on a manifold is *symmetric* provided that

$$\nabla_X Y - \nabla_Y X = [X, Y].$$

Exercise 0.5. Show that a connection is symmetric if and only the corresponding Christoffel symbols satisfy

$$\Gamma_{ij}^k = \Gamma_{ji}^k$$

in every local chart.

If a connection is compatible with the metric and is symmetric we say that it is *Riemannian*. The following result is known as the fundamental theorem of Riemannian Geometry

Theorem 0.6. *Every Riemannian manifold admits a unique Riemannian connection.*

Proof. First suppose that the manifold (M, g) does admit some Riemannian connection ∇ . We will show then that ∇ is unique. To see this, first note that, for any vector fields $X, Y, Z \in \mathcal{M}$,

$$\begin{aligned} Zg(X, Y) &= g(\nabla_Z X, Y) + g(X, \nabla_Z Y), \\ Xg(Y, Z) &= g(\nabla_X Y, Z) + g(Y, \nabla_X Z), \\ Yg(Z, X) &= g(\nabla_Y Z, X) + g(Z, \nabla_Y X). \end{aligned}$$

This yields that

$$\begin{aligned} Zg(X, Y) + Xg(Y, Z) - Yg(Z, X) \\ = g([X, Z], Y) + g([Y, Z], X) + g([X, Y], Z) + 2g(Z, \nabla_Y X). \end{aligned}$$

Therefore

$$g(Z, \nabla_Y X) = \frac{1}{2} \left(Zg(X, Y) + Xg(Y, Z) - Yg(Z, X) \right. \\ \left. - g([X, Z], Y) - g([Y, Z], X) - g([X, Y], Z) \right).$$

This shows that $\nabla_Y X$ is completely determined by g , so it must be unique.

To prove existence, now note that we may define ∇ by using the last expression displayed above. It is easy to check that ∇ would then be a Riemannian connection. \square

Next we are going to derive the local expression for the Christoffel symbols associated to a Riemannian connection. Let (U, ϕ) be a local chart of M and $E_i(p) := d\phi_{\phi(p)}^{-1}(e_i)$ be the corresponding coordinate vector fields on U . Then, recalling the $[E_i, E_j] = 0$, the last displayed expression yields that

$$g\left(E_k, \sum_{\ell} \Gamma_{ij}^{\ell} E_{\ell}\right) = \frac{1}{2} \left(E_k g(E_i, E_j) + E_i g(E_j, E_k) - E_j g(E_k, E_i) \right).$$

Now set $g_{ij} := g(E_i, E_j)$. Further recall that if $f: M \rightarrow \mathbf{R}$ is any function then $E_i f(p) = D_i(f \circ \phi^{-1})(\phi(p))$. Thus if we set $\bar{f} := f \circ \phi^{-1}$, then $E_i(f)$, then we have $E_i f(p) = D_i \bar{f}(\phi(p))$, and the last expression may be rewritten as:

$$\sum_{\ell} \bar{\Gamma}_{ij}^{\ell} \bar{g}_{k\ell} = \frac{1}{2} \left(D_k \bar{g}_{ij} + D_i \bar{g}_{jk} - D_j \bar{g}_{ki} \right).$$

Now let g^{ij} be the coefficients of the matrix which is the inverse of the matrix with coefficients g_{ij} . Then $\sum_k \bar{g}_{k\ell} g^{km} = \delta_{\ell m}$ where $\delta_{\ell m}$ are the coefficients of the identity matrix. Therefore

$$\sum_{k\ell} \bar{\Gamma}_{ij}^{\ell} \bar{g}_{k\ell} \bar{g}^{km} = \sum_{\ell} \bar{\Gamma}_{ij}^{\ell} \delta_{\ell m} = \bar{\Gamma}_{ij}^m.$$

This yields that

$$\bar{\Gamma}_{ij}^m = \frac{1}{2} \sum_k \bar{g}^{km} \left(D_k \bar{g}_{ij} + D_i \bar{g}_{jk} - D_j \bar{g}_{ki} \right). \quad (1)$$

0.3 Induced connection on Riemannian submanifolds

Recall that if \bar{M} is a manifold with connection $\bar{\nabla}$, then any submanifold $M \subset \bar{M}$ inherits a connection ∇ given by

$$\nabla_X Y := (\bar{\nabla}_{\bar{X}} \bar{Y})^{\top}.$$

Further recall that, if (\bar{M}, \bar{g}) is a Riemannian manifold, then M inherits a Riemannian metric g given by

$$g(X, Y) := \bar{g}(\bar{X}, \bar{Y}).$$

Thus one may ask that if $\bar{\nabla}$ is the Riemannian connection of \bar{M} , then is ∇ a Riemannian connection, i.e., is it symmetric and is compatible with g ? Here we show that the answer is yes:

Proposition 0.7. *The induced connection on a Riemannian submanifold is Riemannian.*

Proof. Let $p \in M$, X, Y be vector fields on M , and \bar{X}, \bar{Y} be their extensions to a neighborhood $U \subset \bar{M}$ of p . Then

$$\begin{aligned} Z_p g(X, Y) &= Z_p \bar{g}(\bar{X}, \bar{Y}) \\ &= \bar{g}(\bar{\nabla}_{Z_p} \bar{X}, Y_p) + \bar{g}(X_p, \bar{\nabla}_{Z_p} \bar{Y}) \\ &= \bar{g}((\bar{\nabla}_{Z_p} \bar{X})^\top, Y_p) + \bar{g}(X_p, (\bar{\nabla}_{Z_p} \bar{Y})^\top) \\ &= g(\nabla_{Z_p} X, Y_p) + g(X_p, \nabla_{Z_p} Y) \end{aligned}$$

So ∇ is compatible with g . Next note that

$$\nabla_{X_p} Y - \nabla_{Y_p} X = (\nabla_{\bar{X}_p} \bar{Y})^\top - (\nabla_{\bar{Y}_p} \bar{X})^\top = [\bar{X}, \bar{Y}]_p^\top.$$

But if \bar{f} is any function on \bar{M} and f is its restriction to M , then

$$[\bar{X}, \bar{Y}]_p \bar{f} = X_p(\bar{Y}f) - Y_p(\bar{X}f) = X_p(Yf) - Y_p(Xf) = [X, Y]_p f.$$

Thus

$$[\bar{X}, \bar{Y}]_p^\top = [X, Y]_p^\top = [X, Y]_p.$$

So ∇ is symmetric. □

0.4 Speed of Geodesics

If (M, g) is a Riemannian metric, we say that a curve $c: I \rightarrow M$ is a (Riemannian) geodesic provided that g is a geodesic with respect to the Riemannian connection of M .

Lemma 0.8. *Every Riemannian geodesic $c: I \rightarrow M$ has constant speed, i.e., $g(c'(t), c'(t))$ is constant.*

Proof. Let \bar{c}' be a vector field in a neighborhood U of $c(t_0)$ such that $\bar{c}'(c(t)) = c'(t)$, for all $t \in [t_0 - \epsilon, t_0 + \epsilon]$. Now define $f: U \rightarrow \mathbf{R}$ by $f(p) = g(\bar{c}'(p), \bar{c}'(p))$. Then $g(c'(t), c'(t)) = f(c(t))$, and it follows that

$$\left. \frac{d}{dt} g(c'(t), c'(t)) \right|_{t=t_0} = (f \circ c)'(t_0) = c'(t_0) [g(\bar{c}', \bar{c}')] = 2g(\nabla_{c'(t_0)} \bar{c}', \bar{c}'(t_0)) = 0.$$

Thus $g(c'(t), c'(t))$ is constant. □

If $c: I \rightarrow M$ is any curve, then we say that $\bar{c}: J \rightarrow M$ is a *reparametrization* of c provided that $\bar{c} = c \circ u$ for some diffeomorphism $u: J \rightarrow I$.

Lemma 0.9. *If $c: I \rightarrow M$ is a geodesic, then so is any reparametrization $\bar{c} = c \circ u$, where $u(t) = kt + t_0$ for some constants k and t_0 .*

Proof. The chain rule yields that

$$\bar{c}'(t) = d\bar{c}_t(1) = dc_{u(t)} \circ du_t(1) = dc_{u(t)}(u'(t)) = dc_{u(t)}(k) = kdc_{u(t)}(1) = kc'(u(t)).$$

Consequently,

$$\nabla_{\bar{c}'(t)}\bar{c}' = \nabla_{kc'(kt+t_0)}kc' = k^2\nabla_{c'(kt+t_0)}c' = 0.$$

□

Proposition 0.10. *Let $c: I \rightarrow M$ be a geodesic. Then any reparametrization $\bar{c}: J \rightarrow M$ of c is a geodesic as well, if and only if it has constant speed.*

Proof. If \bar{c} is a geodesic, then it must have constant speed as we showed earlier. Now suppose that \bar{c} has constant speed. Further note that, since $\bar{c} = c \circ u$, for some diffeomorphism $u: J \rightarrow I$, it follows that

$$\bar{c}'(t) = d\bar{c}_t(1) = dc_{u(t)} \circ du_t(1) = dc_{u(t)}(u'(t)) = u'(t)dc_{u(t)}(1) = u'(t)c'(u(t)).$$

Thus, since \bar{c}' and c' both have constant magnitudes, it follows that u' is constant. But then $u(t) = kt + t_0$, and the previous lemma implies that \bar{c} is a geodesic. □

0.5 Example: Geodesics of \mathbf{H}^2

Here we show that the (nontrivial) geodesics in the Poincaré's upper half-plane either trace vertical lines or semicircles which meet the x -axis orthogonally. To this end, we first recall that the standard (hyperbolic) metric on the upper half plane is given by

$$g_{(x,y)}(X, Y) = \frac{\langle X, Y \rangle}{y^2}.$$

Thus

$$g_{11}(x, y) = \frac{1}{y^2}, \quad g_{12}(x, y) = g_{21}(x, y) = 0, \quad g_{22}(x, y) = \frac{1}{y^2}.$$

Further

$$g^{11}(x, y) = y^2, \quad g^{12}(x, y) = g^{21}(x, y) = 0, \quad g^{22}(x, y) = y^2.$$

Now note that we may let the local chart ϕ to be the identity function. Then $\bar{\Gamma}_{ij}^m = \Gamma_{ij}^m$, and so using (1) we may compute that

$$\Gamma_{11}^1 = \Gamma_{12}^2 = \Gamma_{22}^1 = 0, \quad \Gamma_{11}^2 = \frac{1}{y}, \quad \Gamma_{12}^1 = \Gamma_{22}^2 = -\frac{1}{y}.$$

Now recall that $c: I \rightarrow \mathbf{H}^2$ is a geodesic if the following equations are satisfied:

$$\ddot{c}^k(t) + \sum_{ij} \dot{c}^i(t)\dot{c}^j(t)\Gamma_{ij}^k(c(t)) = 0.$$

So if $c(t) = (x(t), y(t))$, then we have

$$\ddot{x} - 2\frac{\dot{x}\dot{y}}{y} = 0, \quad \ddot{y} + \frac{\dot{x}^2 - \dot{y}^2}{y} = 0. \quad (2)$$

To find the solution to these equations, subject to initial conditions $c(0) = (x_0, y_0)$ and $\dot{c}(0) = (\dot{x}_0, \dot{y}_0)$, first suppose that $\dot{x}_0 = 0$. Then the second equation reduces to $\dot{y}/y = \text{const}$. Thus, when $\dot{x}_0 = 0$, then either c traces a vertical line (if $\dot{y}_0 \neq 0$) or is just a point (if $\dot{y}_0 = 0$). It remains then to consider the case when $\dot{x}_0 \neq 0$. We claim that in this case c traces a part of a circle centered at a point on the x -axis, i.e.,

$$(x - a)^2 + y^2 = \text{const}$$

for some constant a (in particular, when $\dot{x}_0 \neq 0$, then $\dot{y}_0 \neq 0$ as well, which may be readily seen from the second equation in (2)). Differentiating both sides of the above equality yields that the above equality holds if and only if

$$a = x + \frac{y\dot{y}}{\dot{x}}.$$

So we just need to check that a is indeed constant, which is a matter of a simple computation with the aid of (2):

$$\dot{a} = \dot{x} + \frac{(\dot{y}^2 + y\ddot{y})\dot{x} - y\dot{y}\ddot{x}}{\dot{x}^2} = \dot{x} + \frac{(\dot{y}^2 + \dot{y}^2 - \dot{x}^2)\dot{x} - \dot{y}(2\dot{x}\dot{y})}{\dot{x}^2} = 0.$$